# Deformation Theory of Asymptotically Conical Spin(7)-Instantons

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#### Abstract

In this talk we discuss the deformation theory of instantons on asymptotically conical Spin(7)-manifolds where the instanton is asymptotic to a fixed nearly  $G_2$ -instanton at infinity. By relating the deformation complex with Dirac operators and spinors, we apply spinorial methods to identify the space of infinitesimal deformations with the kernel of the twisted negative Dirac operator on the asymptotically conical Spin(7)-manifold.

Finally we apply this theory to describe deformations of Fairlie-Nuyts-Fubini-Nicolai (FNFN) Spin(7)-instantons on  $\mathbb{R}^8$ , where  $\mathbb{R}^8$  is considered to be an asymptotically conical Spin(7)-manifold asymptotic to the cone over  $S^7$ . We also calculate the virtual dimension of the moduli space using Atiyah-Patodi-Singer index theorem and the spectrum of the twisted Dirac operator.<sup>1</sup>

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### **1** Preliminaries

#### **1.1** Nearly $G_2$ - and Spin(7)-Manifolds

**Definition 1.1.** Let  $\Sigma$  be a Riemannian 7-dimensional manifold. A 3-form  $\phi \in \Omega^3(\Sigma)$  is called a  $G_2$ -structure on  $\Sigma$  if in local orthonormal frame  $e^1, \ldots, e^7, \phi$  can be written as

$$\phi = e^{127} + e^{347} + e^{567} + e^{145} + e^{136} + e^{235} - e^{246}$$
(1.1)

where  $e^{ijk} := e^i \wedge e^j \wedge e^k$ . The group  $G_2$  is the stabilizer group of  $\phi$  restricted to each tangent space, and is a 14-dimensional simple, connected, simply connected Lie group.

**Definition 1.2.** Let  $\Sigma$  be a 7-dimensional Riemannian manifold and  $\phi \in \Omega^3(\Sigma)$ . Then  $\phi$  is called a *nearly (parallel)*  $G_2$ -structure on  $\Sigma$  if it satisfies

$$d\phi = \tau_0 \psi \tag{1.2}$$

where  $\psi = *\phi$  and  $\tau_0 \in \mathbb{R} \setminus \{0\}$ . In this case,  $(\Sigma, \phi)$  is called a *nearly*  $G_2$ -manifold.

**Definition 1.3.** Let X be an 8-dimensional Riemannian manifold equipped with a 4-form  $\Phi \in \Omega^4(X)$  such that in local orthonormal basis  $e^0, e^1, \ldots, e^7$ , we have  $\Phi = e^0 \wedge \phi + \psi$  where  $\phi$  is as in (1.1) and  $*(e^0 \wedge \phi) = \psi$ . Then  $\Phi$  is said to be a Spin(7)-structure on X and  $(X, \Phi)$  is said to be an almost Spin(7)-manifold.

If  $\Phi$  is torsion-free, i.e., if  $\nabla \Phi = 0$  where  $\nabla$  is the Levi–Civita connection, or equivalently, if  $d\Phi = 0$ , then  $(X, \Phi)$  is called a Spin(7)-manifold.

The 21-dimensional Lie group Spin(7) is the stabiliser group of  $\Phi$  restricted to each tangent space.

#### **1.2** Asymptotically Conical Spin(7)-Manifolds

Let  $(\Sigma, g_{\Sigma})$  be a Riemannian 7-manifold with a nearly  $G_2$ -structure  $\phi$  satisfying  $d\phi = 4\psi$  where  $\psi = *\phi$ . A Spin(7)-cone on  $\Sigma$  is  $C(\Sigma) := (0, \infty) \times \Sigma$  together with a Spin(7)-structure  $(C(\Sigma), \Phi_C)$  defined by

$$\Phi_C := r^3 dr \wedge \phi + r^4 \psi \tag{1.3}$$

where  $r \in (0, \infty)$  is the coordinate.  $\Sigma$  is called the *link* of the cone. The metric  $g_C$  compatible with  $\Phi_C$  is given by

$$g_C = dr^2 + r^2 g_\Sigma \tag{1.4}$$

We note that condition  $d\phi = 4\psi$  implies the torsion free condition  $d\Phi_C = 0$ , which implies that  $(C(\Sigma), g_C, \Phi)$  is a Spin(7)-manifold.

A Spin(7)-cone is not complete. Hence, we consider complete Spin(7)-manifolds whose geometry is asymptotic to the given (incomplete)  $G_2$ -cone.

**Definition 1.4.** Let  $(X, g, \Phi)$  be a non-compact Spin(7)-manifold. X is called an *asymptotically* conical (AC) Spin(7)-manifold with rate  $\nu < 0$  if there exists a compact subset  $K \subset X$ , a compact connected nearly  $G_2$  manifold  $\Sigma$ , and a constant R > 1 together with a diffeomorphism

$$h: (R, \infty) \times \Sigma \to X \setminus K \tag{1.5}$$

such that

$$\left|\nabla_{C}^{j}(h^{*}(\Phi|_{X\setminus K}) - \Phi_{C})\right|(r, p) = O(r^{\nu - j}) \quad \text{as } r \to \infty$$
(1.6)

for each  $p \in \Sigma$ ,  $j \in \mathbb{Z}_{\geq 0}$ ,  $r \in (R, \infty)$ ; where  $\nabla_C$  is the Levi–Civita connection for the cone metric  $g_C$  on  $C(\Sigma)$ , and the norm is induced by the metric  $g_C$ .

 $X \setminus K$  is called the *end* of X and  $\Sigma$  the *asymptotic link* of X.

### 1.3 Lockhart–McOwen Analysis on AC Spin(7)-manifold

Let  $\pi: E \to X$  be a vector bundle over X with a fibre-wise metric and a connection  $\nabla$  compatible with the metric.

**Definition 1.5.** Let  $p \ge 1, k \in \mathbb{Z}_{\ge 0}, \nu \in \mathbb{R}$  and  $C_c^{\infty}(E)$  be the space of compactly supported smooth sections of E. We define the *conically damped* or *weighted Sobolev space*  $W_{\nu}^{k,p}(E)$  of sections of E over X of weight  $\nu$  as follows:

For  $\xi \in C_c^{\infty}(E)$ , we define the weighted Sobolev norm  $\|\cdot\|_{W^{k,p}_{w}(E)}$  as

$$\|\xi\|_{W^{k,p}_{\nu}(E)} = \left(\sum_{j=0}^{k} \int_{X} \left|\varrho^{-\nu+j} \nabla^{j} \xi\right|^{p} \varrho^{-8} \operatorname{dvol}\right)^{1/p}$$
(1.7)

The map  $\rho: X \to \mathbb{R}$ , called a *radius function*, is defined by

$$\varrho(x) := \begin{cases}
1 & \text{if } x \in \text{the compact subset } K \subset X \\
r & \text{if } x = h(r, p) \text{ for some } r \in (2R, \infty), \, p \in \Sigma \\
\widetilde{r} & \text{if } x \in h((R, 2R) \times \Sigma)
\end{cases}$$
(1.8)

where  $h: (R, \infty) \times \Sigma \to X \setminus K$  is the diffeomorphism, and  $\tilde{r}$  is a smooth interpolation between its definition at infinity and its definition on K, in a decreasing manner. Then the weighted Sobolev space  $W_{\nu}^{k,p}(E)$  is the completion of  $C_c^{\infty}(E)$  with respect to the norm  $\|\cdot\|_{W_{\nu}^{k,p}(E)}$ .

We consider  $E := \Lambda^* T^* X \otimes \mathfrak{g}_P$ , and use the notation  $\Omega^{m,k}_{\nu}(\mathfrak{g}_P) := W^{2,k}_{\nu}(\Lambda^m T^* X \otimes \mathfrak{g}_P).$ 

### 2 Asymptotically Conical Spin(7)-Instantons

#### 2.1 Asymptotically Conical Spin(7)-Instantons and Moduli Space

**Definition 2.1.** Let X be an AC Spin(7)-manifold asymptotic to the cone  $C(\Sigma)$ . Let  $P \to X$  be a principal G-bundle over X. P is called asymptotically framed if there exists a principal bundle  $Q \to \Sigma$  such that

$$h^*P \cong \pi^*Q$$

where  $\pi : C(\Sigma) \to \Sigma$  is the natural projection.

Such framing always exists. So we fix a framing Q.

**Definition 2.2.** Let X be an AC Spin(7)-manifold asymptotic to the cone  $C(\Sigma)$ . Let  $P \to X$  be an asymptotically framed bundle. A connection A on P is called an *asymptotically conical* connection with rate  $\nu$  if there exists a connection  $A_{\Sigma}$  on  $Q \to \Sigma$  such that

$$\left|\nabla_C^j(h^*(A) - \pi^*(A_{\Sigma}))\right| = O(r^{\nu - 1 - j}) \quad \text{as } r \to \infty$$
(2.1)

for each  $p \in \Sigma$ ,  $j \in \mathbb{Z}_{\geq 0}$ ,  $\nu < 0$ . The norm is induced by the cone metric and the metric on  $\mathfrak{g}$ .

A is called *asymptotic* to  $A_{\Sigma}$  and  $\nu_0 := \inf\{\nu : A \text{ is AC with rate } \nu\}$  is called the *fastest rate of* convergence of A.

Let  $G \to GL(V)$  be a faithful representation of G, and consider the associated vector bundle End(V). We define the *weighted gauge group* by

$$\mathcal{G}_{k+1,\nu} := \{ \varphi \in C^0(\operatorname{End}(V)) : \|I - \varphi\|_{k+1,\nu} < \infty, \varphi \in G \}$$

We also define  $\mathcal{G}_{\nu} := \bigcap_{l=1}^{\infty} \mathcal{G}_{l,\nu}$ .

A connection A on P is a Spin(7)-instanton if the curvature  $F_A$  satisfies  $*(\Phi \wedge F_A) = -F_A$ . The moduli space of Spin(7)-instantons asymptotic to  $A_{\Sigma}$  with rate  $\nu$  is given by

 $\mathcal{M}(A_{\Sigma},\nu) := \{Spin(7) \text{ instanton } A \text{ on } P \text{ satisfying } (2.1) \text{ asymptotic to } A_{\Sigma}\}/\mathcal{G}_{\nu}$ 

#### 2.2 Deformations of Asymptotically Conical Spin(7)-Instantons

Let A be an asymptotically conical reference connection that also satisfies the Spin(7)-instanton equation. Then, we have  $\pi_7(F_A) = 0$ . Now, we can write any other connection in some open neighbourhood of A as  $A' = A + \alpha$  for  $\alpha \in \Omega^1(\mathfrak{g}_P)$ . Then,

$$F_{A'} - F_A = d_A \alpha + \frac{1}{2} [\alpha, \alpha].$$

Hence the connection A' is a Spin(7)-instanton if and only if  $\pi_7(F_{A+\alpha}) = 0$ , i.e.,

$$\pi_7\left(d_A\alpha + \frac{1}{2}[\alpha,\alpha]\right) = 0.$$

We also have the gauge fixing condition  $d_A^* \alpha = 0$ . We consider the non-linear operator

$$\mathfrak{D}_{A}^{\mathrm{NL}}: \Gamma(\Lambda^{1} \otimes \mathfrak{g}_{P}) \to \Gamma((\Lambda^{0} \oplus \Lambda^{2}_{7}) \otimes \mathfrak{g}_{P})$$
$$\alpha \mapsto \left(d_{A}^{*}\alpha, \pi_{7}\left(d_{A}\alpha + \frac{1}{2}[\alpha, \alpha]\right)\right)$$
(2.2)

Hence, the local moduli space of Spin(7)-instanton can be expressed as the zero set of  $\mathcal{D}_A^{\text{NL}}$ , i.e.,  $\left(\mathcal{D}_A^{\text{NL}}\right)^{-1}(0)$ .

Let  $\Sigma$  be a nearly  $G_2$ -manifold and  $Q \to \Sigma$  is a principal bundle. Let X be an AC Spin(7)manifold with link  $\Sigma$ , and let  $P \to X$  be an asymptotically framed bundle. Let  $A_{\Sigma}$  be a connection on Q. Consider the Dirac operators

$$\mathfrak{P}_{A_{\Sigma}}: \Gamma(\mathfrak{F}(\Sigma) \otimes \mathfrak{g}_{Q}) \to \Gamma(\mathfrak{F}(\Sigma) \otimes \mathfrak{g}_{Q})$$
$$\mathfrak{P}_{A}^{-}: \Gamma(\mathfrak{F}^{-}(X) \otimes \mathfrak{g}_{P}) \to \Gamma(\mathfrak{F}^{+}(X) \otimes \mathfrak{g}_{P})$$

Theorem 2.1. The Dirac operator

$$\mathfrak{D}_{A}^{-}: W_{\nu-1}^{k+1,2}(\mathfrak{F}^{-}(X) \otimes \mathfrak{g}_{P}) \to W_{\nu-2}^{k,2}(\mathfrak{F}^{+}(X) \otimes \mathfrak{g}_{P})$$

is Fredholm if  $\nu$  is not a critical weight, i.e.,  $\nu + \frac{5}{2} \in \mathbb{R} \setminus \operatorname{Spec} \mathfrak{P}_{A_{\Sigma}}$ . Moreover, for two non-critical weights  $\nu, \nu'$  with  $\nu \leq \nu'$ , the jump in the index is given by

$$\operatorname{Index}_{\nu'} \mathfrak{D}_{A}^{-} - \operatorname{Index}_{\nu} \mathfrak{D}_{A}^{-} = \sum_{\nu < \lambda < \nu'} \dim \ker \left( \mathfrak{D}_{A_{\Sigma}} - \lambda - \frac{5}{2} \right).$$

**Definition 2.3.** For  $\nu < 0$  the space of infinitesimal deformations is defined to be

$$\mathcal{I}(A,\nu) := \left\{ \alpha \in \Omega^{1,k+1}_{\nu-1}(\mathfrak{g}_P) : \mathcal{D}_A^- \alpha = 0 \right\}.$$
(2.3)

The obstruction space  $\mathcal{O}(A,\nu)$  is a finite-dimensional subspace of  $\Omega_{\nu-2}^{0,k}(\mathfrak{g}_P) \oplus \Omega_{\nu-2}^{2,k}(\mathfrak{g}_P)$  such that,

$$\Omega^{0,k}_{\nu-2}(\mathfrak{g}_P) \oplus \Omega^{2,k}_{\nu-2}(\mathfrak{g}_P) = \mathcal{D}_A^-\left(\Omega^{1,k+1}_{\nu-1}(\mathfrak{g}_P)\right) \oplus \mathcal{O}(A,\nu).$$
(2.4)

We note that  $\mathcal{I}(A, \nu)$  and  $\mathcal{O}(A, \nu)$  are precisely the kernel and cokernel of the twisted Dirac operator corresponding to the rate  $\nu$ . We have the main theorem:

**Theorem 2.2.** Let A be an AC Spin(7)-instanton asymptotic to a nearly  $G_2$  instanton  $A_{\Sigma}$ . Moreover, let  $\nu \in (\mathbb{R} \setminus \mathscr{D}(\mathfrak{P}_A^-)) \cap (-6,0)$ . Then there exists an open neighbourhood  $\mathcal{U}(A,\nu)$  of 0 in  $\mathcal{I}(A,\nu)$ , and a smooth map  $\kappa : \mathcal{U}(A,\nu) \to \mathcal{O}(A,\nu)$ , with  $\kappa(0) = 0$ , such that an open neighbourhood of  $0 \in \kappa^{-1}(0)$  is homeomorphic to a neighbourhood of A in  $\mathcal{M}(A_{\Sigma},\nu)$ . Hence, the virtual dimension of the moduli space is given by dim  $\mathcal{I}(A,\nu) - \dim \mathcal{O}(A,\nu)$ . Moreover,  $\mathcal{M}(A_{\Sigma},\nu)$  is a smooth manifold if  $\mathcal{O}(A,\nu) = \{0\}$ .

### **3** Deformations of The FNFN Spin(7)-Instanton

Let us consider the asymptotically conical Spin(7)-manifold  $\mathbb{R}^8$  asymptotic to the nearly  $G_2$  manifold  $\Sigma = S^7$ . We consider  $S^7$  as a homogeneous nearly  $G_2$  manifold  $Spin(7)/G_2$ . Then we have the canonical bundle  $G_2 \to Spin(7) \to S^7$  (call this bundle P). Also consider the bundle  $Spin(7) \to Spin(7) \times_{(G_2,\iota)} Spin(7) \to S^7$  (call this bundle Q) where  $\iota : G_2 \hookrightarrow Spin(7)$  is the inclusion. This bundle is (bundle) isomorphic to the trivial bundle  $Spin(7) \to Spin(7) \times S^7 \to S^7$ .

Let  $A_{\text{flat}}$  be a Spin(7)-invariant flat connection given by  $A_{\text{flat}} = A_{\Sigma} + a$ . Let  $(r, \sigma) \in (0, \infty) \times S^7$ . Then the connection  $A(r, \sigma) = A_{\Sigma}(\sigma) + f(r)a(\sigma)$  where  $f(r) = \frac{1}{Cr^2 + 1}$  for C > 0 is a function on  $\mathbb{R}^8$ , is an instanton on  $\mathbb{R}^8$ , shall be called the *FNFN Spin*(7)-*instanton*. Clearly the FNFN Spin(7)-instanton A is asymptotic to the canonical connection  $A_{\Sigma}$  with fastest rate of convergence -2.

#### 3.1 The Main Result

**Theorem 3.1.** The virtual dimension of the moduli space  $\mathcal{M}(A_{\Sigma}, \nu)$  of the FNFN Spin(7)-instantons with decay rate  $\nu \in (-2, 0) \setminus \{-1\}$  is given by

virtual-dim 
$$\mathcal{M}(A_{\Sigma}, \nu) = \begin{cases} 1 & \text{if } \nu \in (-2, -1) \\ 9 & \text{if } \nu \in (-1, 0). \end{cases}$$

The dimensional jump of the moduli space happens for the rate -1, which corresponds to the eigenvalue 3/2.

From the fact that the eigenvalues of the twisted Dirac operator in the range [1/2, 5/2] are 1/2and 3/2, corresponding to the trivial and spin representations respectively, we should expect that the rate of dilation should be 1/2 - 5/2 = -2 and that of translation should be 3/2 - 5/2 = -1. This can be easily verified from the fact that the two deformations translation and dilation are given by  $\iota_{\frac{\partial}{\partial r^i}} F_A$  and  $\iota_{x^i \frac{\partial}{\partial r^i}} F_A$  respectively.

#### **3.2** Twisted Dirac Operators on $Spin(7)/G_2$

We start with a reductive homogeneous bundle G/H. Hence  $\mathfrak{g} = \mathfrak{m} \oplus \mathfrak{h}$ . Let  $\{I_A\}$  is an orthonormal basis of  $\mathfrak{g}$ . We want to figure out the spectrum of the twisted Dirac operator

$$\mathfrak{P}_{A_{\Sigma}}: \Gamma(\mathfrak{f}_{\mathbb{C}}(\Sigma) \otimes (\mathfrak{g}_{P})_{\mathbb{C}}) \to \Gamma(\mathfrak{f}_{\mathbb{C}}(\Sigma) \otimes (\mathfrak{g}_{P})_{\mathbb{C}})$$

Using Frobenius reciprocity, we decompose the spinor bundle as

$$L^{2}(\mathscr{S}_{\mathbb{C}}(\Sigma)\otimes(\mathfrak{g}_{P})_{\mathbb{C}})\cong L^{2}(G,\Delta\otimes V)^{H}\cong\bigoplus_{\gamma\in\widehat{G}}\operatorname{Hom}(V_{\gamma},\Delta\otimes V)^{H}\otimes V_{\gamma}$$
(3.1)

where  $V_{(a,b,c)}$  is an irreducible representation of  $\mathfrak{spin}(7)$  and  $V_{(a,b)}$  is an irreducible representation of  $\mathfrak{g}_2$ .

#### 3.3 Eigenvalues of the Twisted Dirac Operator

For an FNFN Spin(7)-instanton, the fastest rate of convergence is -2. Hence we consider the family of moduli spaces  $\mathcal{M}(A_{\Sigma},\nu)$  for  $\nu \in (-2,0)$ . We want to find the critical weights in (-2,0), i.e.,  $\nu \in (-2,0)$  such that  $\nu + \frac{5}{2} \in \operatorname{Spec} \mathfrak{P}_{A_{\Sigma}}$ . Hence, we are interested in finding all the eigenvalues of the twisted Dirac operator in the interval  $\left[-\frac{5}{2}, \frac{5}{2}\right]$ . By an eigenvalue bound calculation, we find that if  $V_{\gamma}$  is not one of the irreducible representations  $V_{(0,0,0)}, V_{(1,0,0)}, V_{(0,1,0)}, V_{(2,0,0)}, V_{(1,0,1)}$ , then the operator  $(\mathfrak{P}_{A_{\Sigma}})_{\gamma}$  has no eigenvalues in the interval  $\left(\frac{1}{2}, \frac{5}{2}\right)$ .

**Proposition 3.2.** By explicit computation, we find that the eigenvalues of the twisted Dirac operator  $\left(\mathfrak{P}_{A_{\Sigma}}^{0}\right)_{\gamma}$  in the interval  $\left[-\frac{5}{2},\frac{5}{2}\right]$  are  $-\frac{5}{2},-\frac{3}{2},\frac{1}{2},\frac{3}{2}$ , and that of in the interval  $\left(\frac{1}{2},\frac{5}{2}\right)$  is  $\frac{3}{2}$  corresponding to the spin representation  $V_{(0,0,1)}$ .

### 3.4 Index of the Twisted Dirac Operator

Let  $g_C$  be the asymptotically conical metric on  $\mathbb{R}^8$ . We define the metric  $g_{CI} := \frac{1}{\rho^2} g_C$ . Then  $M := (\mathbb{R}^8, g_{CI})$  resembles a cigar (the reason  $g_{CI}$  is usually called a *cigar metric*).



Then,

**Proposition 3.3.** Let  $B_R^8 := \{x \in \mathbb{R}^8 : |x| \leq R\}$  be 8-dimensional ball of radius R. Then for sufficiently large R, we have

Index 
$$\left( \mathfrak{P}_{A,CI}^{-}, \mathbb{R}^{8}, g_{CI} \right)$$
 = Index  $\left( \mathfrak{P}_{\widetilde{A},CI}^{-}, B_{R}^{8}, g_{CI} \right)$ .

Moreover, for sufficiently large T, we have

$$\operatorname{Index}\left(\boldsymbol{\mathfrak{P}}_{A}^{-}, \mathbb{R} \times S^{7}, g\right) = \operatorname{Index}\left(\boldsymbol{\mathfrak{P}}_{\widetilde{A}'}^{-}, [-T, T] \times S^{7}, g\right)$$

where g is the cylindrical metric  $g = dt^2 + g_{S^7}$ .

Using Atiyah–Patodi–Singer index theorem us find that the index of the Dirac operator  $\mathcal{P}_{A_{CI}}^{-}$ on  $\mathscr{S}(\mathbb{R}^8, g_{CI})$  twisted by the bundle  $\mathfrak{spin}(7)$  is

$$\operatorname{Ind} \mathfrak{D}_{A_{CI}}^{-} = I(M) + CS(S^{7}) + \frac{1}{2}\eta(S^{7})$$
$$= -\frac{1}{12} \int_{\mathbb{R}\times S^{7}} \left( p_{1}(\mathfrak{g}_{P})^{2} - 2p_{2}(\mathfrak{g}_{P}) \right) + 0$$
$$+ \frac{1}{12} \int_{\mathbb{R}\times S^{7}} \left( p_{1}(\mathfrak{g}_{P})^{2} - 2p_{2}(\mathfrak{g}_{P}) \right)$$
$$= 0.$$

Hence,

$$\operatorname{Ind}_{-\frac{5}{2}} \mathfrak{D}_A^- = \operatorname{Ind} \mathfrak{D}_{A_{CI}}^- = 0.$$

## References

- [1] J. Driscoll. Deformations of asymptotically conical  $G_2$  instantons. arXiv:1911.01991v3, 2021.
- [2] D. B. Fairlie and J. Nuyts. Spherically symmetric solutions of gauge theories in eight dimensions. J. Phys. A: Math. Gen., 17(14):2867–2872, 1984.
- [3] S. Fubini and H. Nicolai. The Octonionic Instanton. Physics Letters B, 155(5-6):369–372, 1985.

- [4] T. Ghosh. Deformation Theory of Asymptotically Conical Spin(7)-Instantons. arXiv:2305.03646, 2023.
- [5] R.B. Lockhart and R.C. McOwen. Elliptic differential operators on noncompact manifolds. Annali della Scuola Normale Superiore di Pisa-Classe di Scienze, 12(3):409–447, 1985.
- [6] S.P. Marshall. Deformations of special Lagrangian submanifolds. PhD thesis, University of Oxford, 2002.
- [7] R. Singhal. Deformations of  $G_2$ -instantons on nearly  $G_2$  manifolds. Ann. Glob. Anal. Geom., 62:329–366, 2022.