SU(3)-structures of LT type

Dimitrios Tsimpis Pau, June 2023









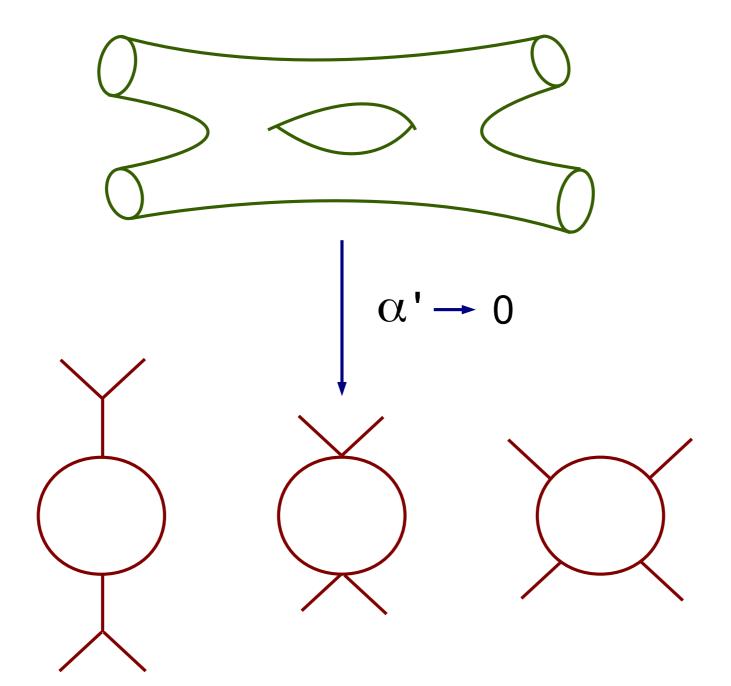
Based on:

- * D. Lüst & DT, hep-th/0412250 (2004)
- * P. Koerber & DT, 0804. 0614 (2008)
- * M. Larfors, D. Lüst & DT, 1005.2194 (2010)
- * *DT*, 1206.5900 (2012)
- * R. Terrisse & DT, 1707.04636 (2017)
- * M. Larfors, A. Lukas, F. Ruehle & DT, 2022

Outline

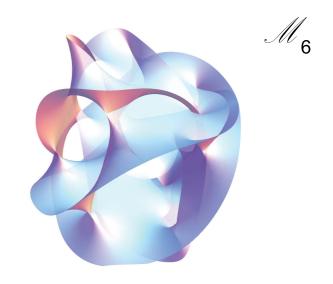
- Introduction: context & motivation
- Characterization of LT solutions
- Constructions of LT solutions
- Conclusions: outlook & wishlist

- String/M-theory is a leading candidate for a consistent theory of quantum gravity.
- At low energies (critical) string theory reduces to 10d supergravity



- A great part of our knowledge of the theory (dualities, holography, branes, black holes,...) comes from the study of its supersymmetric, (bosonic) supergravity vacua (solutions)
- Iod supergravity contains (in addition to the graviton) fermions and higher-rank antisymmetric tensors (flux).

- In the absence of flux, *susy* vacua of the theory are of the form $\mathbb{R}^{1,3} \times \mathcal{M}_6$ where \mathcal{M}_6 is a Calabi-Yau manifold.
- Good control of the physics, beautiful mathematics.
- Powerful tools from math.AG.
- * Candelas, Horowitz, Strominger, Witten, 1985
- * Strominger, Witten, 1985
- * http://hep.itp.tuwien.ac.at/-kreuzer/CY/



- Good physical reasons to turn on the flux (generic case).
- A famous no-go theorem excludes Minkowski *flux* vacua, provided:
 - absence of sources, no (or mild) singularities
 - compactness
 - two-derivative actions
 - the Strong Energy Condition is obeyed by the 10d/11d theory
- * Gibbons, 1984
- * Maldacena & Nuñez, 2000

- In the presence of flux: *susy* vacua of the form $AdS_4 \times \mathcal{M}_6$ where \mathcal{M}_6 is *not* special-holonomy.
- Holography, moduli stabilisation, de Sitter uplifts, ...
- * Freund, Rubin, 1980
- * Duff, Pope, 1982
- * Nilsson, Pope, 1984

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- * Freund, Rubin, 1980
- * Duff, Pope, 1982
- * Nilsson, Pope, 1984
- G-structures, generalized (complex) geometry, ...
- * Gauntlett, Kim, Martelli, Waldram, 2001
- * Graña, Minasian, Petrini, Tomasiello, 2004

- LT-vacua are $\mathcal{N} = 1$ solutions of (massive) 10d IIA supergravity of the form $AdS_4 \times \mathcal{M}_6$ with \mathcal{M}_6 in a certain subclass of *half-flat*.
- * D. Lüst & DT, 2004

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- * De Wolfe, Giryavets, Kachru, Taylor, 2005
- * Acharya, Benini, Valandro, 2007

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- * Acharya, Benini, Valandro, 2007
- Relevant to the question of scale separation and the swampland
- * Many recent papers

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IIA supergravity

IIA (bosonic) action

$$S = \frac{1}{2\kappa_{10}^2} \int d^{10}x \sqrt{g} \left(-R + \frac{1}{2} (\partial \phi)^2 + \frac{1}{2 \cdot 2!} e^{3\phi/2} F^2 + \frac{1}{2 \cdot 3!} e^{-\phi} H^2 + \frac{1}{2 \cdot 4!} e^{\phi/2} G^2 + \frac{1}{2} m^2 e^{5\phi/2} \right) + S_{CS}$$

Bianchi identities

$$dF = mH$$
; $dH = 0$; $dG = H \wedge F$

Supersymmetry of the (bosonic) vacuum

$$\nabla_{M} \epsilon = \frac{1}{16} m e^{5\phi/4} \Gamma_{M} \epsilon + \frac{1}{32} e^{3\phi/4} \Gamma_{MNP} F^{NP} \epsilon + \dots$$
$$0 = \left(10 m e^{5\phi/4} - 3 e^{3\phi/4} \Gamma_{MN} F^{MN} + \dots \right) \epsilon$$

- An SU(3)-structure on \mathcal{M}_6
 - \blacksquare A complex decomposable 3-form Ω
 - A real 2-form J such that

$$\Omega \wedge J = 0 \; ; \quad \Omega \wedge \Omega^* = \frac{4i}{3}J \wedge J \wedge J \neq 0$$

- Equivalences:
 - An SU(3)-structure on \mathcal{M}_6
 - Existence on \mathcal{M}_6 of (g, \mathcal{I}) with $c_1(\mathcal{I}) = 0$
 - Existence on \mathcal{M}_6 of (g, η) with η pure & nowhere-vanishing

- Link with supergravity:
 - Susy parameter: $\epsilon \sim \zeta \otimes \eta$
 - SU(3)-structure: $\Omega \sim \eta \gamma_{(3)} \eta$; $J \sim \eta^{\dagger} \gamma_{(2)} \eta$
- Torsion classes:

$$dJ = \frac{3}{2} \operatorname{Im}(W_1^* \Omega) + W_4 \wedge J + W_3$$
$$d\Omega = W_1 J \wedge J + W_2 \wedge J + \Omega \wedge W_5^*$$

where:

$$W_1 \sim 1 \oplus 1; \quad W_2 \sim 8 \oplus 8; \quad W_3 \sim 6 \oplus \overline{6}; \quad W_4 \sim 3 \oplus \overline{3}; \quad W_5 \sim 3$$

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 - Susy parameter: $\epsilon \sim \zeta \otimes \eta$
 - SU(3)-structure: $\Omega \sim \eta \gamma_{(3)} \eta$; $J \sim \eta^{\dagger} \gamma_{(2)} \eta$
- Torsion classes:

$$\nabla_m \eta = \frac{1}{2} \left(W_{4m}^{(1,0)} + W_{5m} - \text{c.c.} \right) \eta$$

$$+\frac{1}{16} \left(4W_1 g_{mn} - 2W_4^p \Omega_{pmn} + 4iW_{2mn} - iW_{3mpq} \Omega^{pq}_n\right) \gamma^n \eta^c$$

where:

$$W_1 \sim 1 \oplus 1; \quad W_2 \sim 8 \oplus 8; \quad W_3 \sim 6 \oplus \overline{6}; \quad W_4 \sim 3 \oplus \overline{3}; \quad W_5 \sim 3$$

- Link with supergravity:
 - Susy equations:

$$\nabla \eta = F \cdot \gamma \eta$$

Compare with:

$$\nabla \eta = W \cdot \gamma \eta$$

- Conclusion:
 - lacktriangle The flux is expressed in terms of torsion classes: $F \sim W$
 - Constraints on the torsion classes.

What about the equations of motion?

An *integrability theorem* guarantees that (potentially in the presence of *calibrated* sources) the Einstein, dilaton and 3-form eom's automatically follow from *susy*, the (generalized) *Bianchi identities* and a certain (mild) assumption on the form of the solution.

- * D. Lüst & DT, 2004
- * P. Koerber & DT, 2007
- * D. Lüst, F. Marchesano, L. Martucci & DT, 2008
- * D. Prins & DT, 2013

The LT solutions

- \blacksquare The rod metric is of the form $AdS_4 \times \mathcal{M}_6$
- An SU(3)-structure on \mathcal{M}_6 with torsion classes:

$$dJ = -\frac{3}{2}iW_1^- \text{Re}\Omega \; ; \quad d\Omega = W_1^- J \wedge J + W_2^- \wedge J$$

Fluxes:

$$H = \frac{2m}{5}e^{\phi} \operatorname{Re}\Omega$$

$$F_2 = \frac{1}{4}ie^{-\phi}W_1^- J + ie^{-\phi}W_2^-$$

$$F_4 = \frac{9}{4}ie^{-\phi}W_1^- \operatorname{vol}_4 + \frac{3m}{10}J \wedge J$$

$$\frac{1}{L^2} = \left(\frac{1}{25}m^2 + \frac{9}{16}|W_1^-|^2\right)e^{2\phi}$$

 \blacksquare Geometrical problem: construct \mathcal{M}_6 with this structure!

The LT solutions

Generically there are D6/O6 sources:

$$dF_2 + mH = j_6$$

$$j_6 \equiv \left(\frac{2m^2}{5}e^{\phi} - \frac{3}{8}|W_1^-|^2e^{-\phi}\right)\text{Re}\Omega + ie^{-\phi}dW_2^-$$

Absence of sources equivalent to:

$$dW_2^- \propto \text{Re}\Omega$$
; $3|W_1^-|^2 - |W_2^-|^2 \ge 0$

The LT solutions

■ The *Nearly Kähler* limit:

$$dJ = -\frac{3}{2}iW_1^- \text{Re}\Omega ; \quad d\Omega = W_1^- J \wedge J$$

- solutions without sources always possible
- four homogeneous spaces: $S^3 \times S^3$, \mathbb{CP}^3 , $\mathrm{Tw}(\mathbb{CP}^2)$, S^6
- * K. Behrndt, M. Cvetic, 2004
- The Nearly Calabi-Yau limit:

$$dJ = 0 ; \quad d\Omega = W_2^- \wedge J$$

solutions without sources never possible

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■ \mathcal{M}_6 of the form G/H with H subgroup of SU(3)

G	H		
G_2	SU(3)		
$SU(3)\times SU(2)^2$	SU(3)		
$\operatorname{Sp}(2)$	$S(U(2)\times U(1))$		
$SU(3)\times U(1)^2$	$S(U(2)\times U(1))$		
$SU(2)^3 \times U(1)$	$S(U(2)\times U(1))$		
SU(3)	$U(1)\times U(1)$		
$SU(2)^2 \times U(1)^2$	$U(1)\times U(1)$		
$SU(3) \times U(1)$	SU(2)		
$\mathrm{SU}(2)^3$	SU(2)		
$SU(2)^2 \times U(1)$	U(1)		
$\mathrm{SU}(2)^2$	1		

Given a basis of the algebras and a coset representative L

$$\{\mathcal{H}_a\}, \ a = 1, \dots, \dim(H) \ ; \ \{\mathcal{K}_i\}, \ i = 1, \dots, \dim(G) - \dim(H)$$

a coframe e^i on G/H is defined

$$L^{-1}dL = e^i \mathcal{K}_i + \omega^a \mathcal{H}_a$$

Ap-form $\phi = \frac{1}{p!}\phi_{i_1...i_p}e^{i_1}\wedge\cdots\wedge e^{i_p}$ is left-invariant iff $f^j{}_{a[i_1}\phi_{i_2...i_p]j} = 0$; $\phi_{i_1...i_p} = \text{const.}$

where
$$[\mathcal{H}_a, \mathcal{H}_b] = f^c{}_{ab}\mathcal{H}_c$$
,
 $[\mathcal{H}_a, \mathcal{K}_i] = f^j{}_{ai}\mathcal{K}_j$,
 $[\mathcal{K}_i, \mathcal{K}_j] = f^k{}_{ij}\mathcal{K}_k + f^a{}_{ij}\mathcal{H}_a$.

■ The exterior differential of a left-invariant form is left-invariant

- lacksquare Construct the most general left-invariant (J,Ω) for each \mathcal{M}_6
- \blacksquare Calculate $(dJ, d\Omega)$
- $\blacksquare \text{ Impose } \tau \subset W_1^- \oplus W_2^-$

	SU(2):	$\times SU(2)$	$\frac{\mathrm{SU}(3)}{\mathrm{U}(1)\times\mathrm{U}(1)}$	$\frac{\operatorname{Sp}(2)}{\operatorname{S}(\operatorname{U}(2) \times \operatorname{U}(1))}$	$\frac{G_2}{SU(3)}$	$\frac{\mathrm{SU}(3) \times \mathrm{U}(1)}{\mathrm{SU}(2)}$
# of parameters	2	4	4	3	2	4
$\mathcal{W}_2^- \neq 0$	No	Yes	Yes	Yes	No	Yes
$j^6 \propto { m Re}\Omega$	Yes	No	Yes	Yes	Yes	No

Topologies: $S^3 \times S^3$, $\operatorname{Tw}(\mathbb{CP}^2)$, \mathbb{CP}^3 , S^6 , $S^5 \times S^1$

* D. Lüst, P. Koerber & DT, 2008

- lacksquare Construct the most general left-invariant (J,Ω) for each \mathcal{M}_6
- \blacksquare Calculate $(dJ, d\Omega)$
- $\blacksquare \text{ Impose } \tau \subset W_1^- \oplus W_2^-$
- Impose $dW_2^- \propto \text{Re}\Omega$; $3|W_1^-|^2 |W_2^-|^2 \ge 0$

	$\mathrm{SU}(2){ imes}\mathrm{SU}(2)$	$\frac{\mathrm{SU}(3)}{\mathrm{U}(1)\times\mathrm{U}(1)}$	$\frac{\operatorname{Sp}(2)}{\operatorname{S}(\operatorname{U}(2) \times \operatorname{U}(1))}$	$\frac{G_2}{SU(3)}$
# of parameters	1	3	2	1
$\mathcal{W}_2^- \neq 0$	No	Yes	Yes	No

* D. Lüst, P. Koerber & DT, 2008

Explicit examples: nilmanifolds

- No solutions in the absence of sources
- \blacksquare Two solutions with sources: on \mathbb{T}^6 and the Iwasawa manifold
- Related solution: \mathbb{T}^2 over K_3
- * P. Koerber & DT, 2008
- * Caviezel, Koerber, Kors, Lüst, DT & Zagermann, 2008

Explicit examples: twistor spaces

A two-parameter solution on

$$\operatorname{Tw}(\mathbb{CP}^2) \cong \frac{SU(3)}{U(1) \times U(1)}$$

- misses one parameter with respect to the coset description
- A two-parameter solution on

$$\operatorname{Tw}(S^4) \cong \mathbb{CP}^3 \cong \frac{Sp(2)}{S(U(2) \times U(1))}$$

- * A. Tomasiello, 2007
- * Feng Xu, 2006

Explicit examples: sphere bundles

- LT vacua on $S^2(B_4)$ with B_4 positive Kähler-Einstein
 - Local SU(2) structure on B_4

$$\hat{\omega} \wedge \hat{\omega}^* = 2\hat{j} \wedge \hat{j} ; \quad \hat{j} \wedge \hat{\omega} = 0 ;$$

$$d\mathcal{P} = 6 \hat{j} ; \quad d\hat{j} = 0 ; \quad d\hat{\omega} = i\mathcal{P} \wedge \hat{\omega}$$

• Global SU(3) structure

$$J = |h|^2 j + \frac{i}{2} K \wedge K^* \; ; \quad \Omega = h^2 \omega \wedge K$$

where

$$j := \cos \theta \, \hat{j} + \sin \theta \, \Re(e^{i\psi} \hat{\omega})$$

$$\omega := -\sin \theta \, \hat{j} + \cos \theta \, \Re(e^{i\psi} \hat{\omega}) + i \, \Im(e^{i\psi} \hat{\omega})$$

$$K := f d\theta + ig(d\psi + \mathcal{P})$$

* R. Terrisse & DT, 2017

(can we eat our cake and have it too?)

- Susy selects a non-integrable complex structure on \mathcal{M}_6 but \mathcal{M}_6 may also admit another, integrable complex structure.
- Use the underlying algebra-geometric description of \mathcal{M}_6 Example: 3d smooth, compact toric varieties
- * M. Larfors, D. Lüst & DT, 2010

(can we eat our cake and have it too?)

- \blacksquare A *d*-dimensional SCTV corresponds to a fan Σ
- A fan is a (certain) collection of (certain) cones generated by

$$G(\Sigma) = \{v_1, \dots v_n\} ; \quad v_i \in N \cong \mathbb{Z}^d$$

- Classification of 2d SCTV
- Partial classification of (minimal) 3d SCTV
 - $\blacksquare \mathbb{CP}^2$ bundles over \mathbb{CP}^1
 - \blacksquare \mathbb{CP}^1 bundles over 2d SCTV
 - lacksquare complete results for $n \leq 8$
- * Miyake & Oda; Oda, 1978
- Some work needed to read off $G(\Sigma)$

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A SCTV can also be described as a symplectic quotient

$$\mathcal{M}_{2d} = \mu^{-1}(0)/U(1)^s$$

Moment maps

$$\mu^{a} \equiv \sum_{i=1}^{n} Q_{i}^{a} |z^{i}|^{2} - \xi^{a}$$

$$a = 1, \dots s ; \quad (z^{1}, \dots, z^{n}) \in \mathbb{C}^{n} ; \quad d = n - s$$

 $\blacksquare U(1)^s$ action on \mathbb{C}^n

$$z^i \longrightarrow e^{i\varphi_a Q^a_i} z^i$$

• Unique topology for $\xi^a \in \mathcal{K}_{\mathcal{M}}$

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• Given $G(\Sigma)$ we determine the $U(1)^s$ charges by solving

$$\sum_{i=1}^{n} Q_i^a v_i = 0$$

 $lacksquare{1}{2}$ Forms Φ on \mathbb{C}^n that are gauge-invariant and vertical

$$\mathcal{L}_{\mathrm{Im}V^a}\Phi=0$$
; $\iota_{V^a}\Phi=\iota_{\bar{V}^a}\Phi=0$ where $V^a\equiv\sum_iQ_i^az^i\partial_{z_i}$ descend to well-defined forms on \mathcal{M}_{2d}

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Sufficient conditions for global SU(3) structure on SCTV (1,0)-form K on \mathbb{C}^n such that

1.
$$P(K) = K$$

$$2. \quad Q^a(K) = \frac{1}{2}Q^a(\Omega_{\mathbb{C}})$$

3.
$$|K|^2 = 2$$

Local SU(2) structure on SCTV

$$\omega = -\frac{i}{2} K^* \cdot \widetilde{\Omega} \; ; \quad j = \widetilde{J} - \frac{i}{2} K \wedge K^*$$

where

$$\widetilde{\Omega} \propto \prod_{a=1}^{\sigma} \iota_{V^a} \Omega_{\mathbb{C}} \; ; \quad \widetilde{J} = P(J_{\mathbb{C}})$$

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■ Global SU(3) structure on SCTV

$$J = \alpha j - \frac{i\beta^2}{2} K \wedge K^* \; ; \quad \Omega = \alpha \beta e^{i\gamma} K^* \wedge \omega$$

- Torsion classes
 - Must be computed case-by-case
 - Generally $W_i \neq 0$
 - Special points with $W_1, W_3, W_4 = 0$
 - lacksquare Exception: the LT vacuum on \mathbb{CP}^3

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- Formalism applicable to non-compact toric varieties
- * Chen, Dasgupta, Franche, Katz, Tatar, 2010
- \blacksquare Many further examples of K constructed
- * M. Larfors, 2013
- Potentially relaxing the non-vanishing condition on K
- * S. Dabholkar, 2013

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- What about the LT vacuum on $Tw(\mathbb{CP}^2)$?
- Modification of the prescription to obtain a global SU(3) structure on (toric) \mathbb{CP}^1 bundles over arbitrary 2d SCTV

i=1

■ Toric 3d U(I) charges

$$Q_I^A = \begin{pmatrix} q_i^a & -n^a & 0\\ 0 & 1 & 1 \end{pmatrix} ; \quad n_a \in \mathbb{N}$$

where q_i^a are the toric 2d U(I) charges

- Prescription works for $n^a = \sum q_i^a$
- Generic torsion classes

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- \blacksquare *SU*(3) structures on CICY from Machine Learning
- * Anderson, Gerdes, Gray, Krippendorf, Raghuram, Ruehle, 2021 proposed the SU(3) ansatz

$$J = \sum_{i=1}^{m} a_i J_i \; ; \quad \Omega = A_1 \Omega_0 + A_2 \Omega_0^*$$

defined on
$$\begin{bmatrix} \mathbb{CP}^{n_1} & q_1^1 & \dots & q_K^1 \\ \vdots & \vdots & \vdots & \vdots \\ \mathbb{CP}^{n_m} & q_1^m & \dots & q_K^m \end{bmatrix} \; ; \qquad \sum_r q_r^i = n_i + 1$$

subject to

$$|A_1|^2 + |A_2|^2 = \sum_{i,j,k=1}^m \Lambda_{ijk} a_i a_j a_k$$

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- *SU*(3) structures on CICY from Machine Learning
- * Anderson, Gerdes, Gray, Krippendorf, Raghuram, Ruehle, 2021 where the are Λ_{ijk} are read off of

$$J_i \wedge J_j \wedge J_k = \frac{3}{4} i \Lambda_{ijk} \Omega_0 \wedge \Omega_0^*$$

■ The ansatz would produce LT SU(3) structures for

$$a_i = \text{const.}$$
; $A_2 = -A_1^* + \text{const.}$

- Unfortunately the ansatz of *Anderson et al* does not satisfy complex decomposability of Ω .
- * M. Larfors, A. Lukas, F. Ruehle & DT, 2022

Conclusions—wishlist

- LT structures are well-motivated in supergravity/string theory
- Still few known explicit examples
- \blacksquare *SU*(3) decomposable ansätze on CYs and/or toric varieties?
- Interplay between *math.AG* and *math.DG*
- Connection with Hitchin uplift to G2. Use 7d technology?
- Existence theorems? Deformations?
- Do all LT-structure manifolds without sources admit a NK limit?
- Ground for new discoveries!