# Cohomogeneity one examples of deformed instantons 

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Unicamp
Workshop BRIDGES: Special geometries and gauge theories

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## Plan of talk

Words in the title:

- deformed instantons: special connections on complex line bundle $L \rightarrow M$ : where $M$ is either Kähler, $G_{2}$ or $\operatorname{Spin}(7)$ manifold
- Nomenclature: Deformed HYM connections, deformed $G_{2} / \operatorname{Spin}(7)$ instantons or deformed Donaldon-Thomas connections
- Cohomogeneity one: $M / G$ is 1-dimensional
- Background on $G_{2}$ geometry
- What are deformed instantons? Where they come from?
- On cohomogeneity one gauge theory
- Examples of $\mathrm{d} G_{2}$-instantons
- Examples of $\mathrm{d} \operatorname{Spin}(7)$-instantons and dHYM connections


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## $G_{2}$-structures on 7-manifolds

- A $G_{2}$-structure on $M^{7}$ is the data of a 3-form $\varphi$ such that at each $p \in M$ $\exists\left\{x_{i}\right\}_{i=1}^{7}$ such that

$$
\begin{aligned}
\left.\varphi\right|_{p}= & d x_{123}+d x_{145}+d x_{167}+d x_{246}-d x_{257}-d x_{347}-d x_{356} \\
= & d x_{1} \wedge\left(d x_{23}+d x_{45}+d x_{67}\right) \\
& +\operatorname{Re}\left(\left(d x_{2}+i d x_{3}\right) \wedge\left(d x_{4}+i d x_{5}\right) \wedge\left(d x_{6}+i d x_{7}\right)\right)
\end{aligned}
$$

- $G_{2} \subset S O(7) \Rightarrow \varphi$ determines a metric and orientation. Explicitly by

$$
\left.\left.\frac{1}{6}(X\lrcorner \varphi\right) \wedge(Y\lrcorner \varphi\right) \wedge \varphi=g(X, Y) \text { vol }
$$

- $\left(M^{7}, \varphi\right)$ "is" $G_{2}$ manifold $\Leftrightarrow \nabla \varphi=0 \Leftrightarrow d \varphi=0=d * \varphi$ (49 PDEs!)
- Trivial example: If ( $P^{6}, h, \omega, \Omega$ ) is $C Y$ then $M^{7}=S_{t}^{1} \times P^{6}$ is $G_{2}$ with


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## Why care about $G_{2}$-manifolds?

(1) Admits a parallel spinor $\Rightarrow \operatorname{Ric}(g)=0: M^{7}$ is an Einstein manifold
(2) Calibrated geometry

## Definition

A $k$-form $\eta$ is a calibration on $(M, g)$ if $d \eta=0$ and $\eta\left(e_{1}, \ldots e_{k}\right) \leq 1 \forall e_{i}$ st $\mid e_{i} \|_{g}=1$

If $L^{k} \subset M$ st $\left.\eta\right|_{L}=\operatorname{vol}_{L}$ then $L$ is calibrated submanifold. Suppose $L$ is cpt calibrated and $\partial N^{k+1}=L \cup \overline{L^{\prime}}$ then

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\operatorname{Vol}(L)=\int_{L} \eta=\int_{L^{\prime}} \eta \leq \operatorname{Vol}\left(L^{\prime}\right) .
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$L$ is minimal $k$-submanifold in homology class! And $L$ is minimiser! On $\left(M^{7}, \varphi\right) G_{2}$-manifold: calibrated by $\varphi(* \varphi)$ means $L^{3(4)}$ is (co)-associative. On $\left(M^{2 n}, \omega, \Omega\right) C Y$-fold: calibrated by $\omega^{k} / k$ ! means $L^{2 k}$ complex sub-mfd, calibrated by $\operatorname{Re}\left(e^{i \theta} \Omega\right)$ means $L$ is special Lagrangian

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## Gauge theory: $G_{2}$-instanton

(3) On $\left(M^{4}, g\right.$, vol), let $A$ be a connection on a vector bundle $E \rightarrow M$ then $F_{A}:=d A+\frac{1}{2}[A \wedge A] \in \Omega^{2}(E n d(E)) \cong \Omega_{+}^{2} \oplus \Omega_{-}^{2}$. Call $F_{A}$ an ASD instanton if $F_{A} \in \Omega_{-}^{2}$ i.e. $* F_{A}=-F_{A}$.


Rmk: For a CY 3-fold $P^{6}$, instanton means (traceless) HYM i.e. $*\left(F_{A} \wedge \omega\right)=-F_{A} \Leftrightarrow F_{A} \wedge \operatorname{Im}(\Omega)=0=F_{A} \wedge \omega \wedge \omega \Leftrightarrow F_{A} \in \mathfrak{s u}(3)=\Omega_{0}^{1}$ Prop. On $M^{7}=S_{t}^{1} \times P^{6}$ as before, $A$ is traceless HYM on $P^{6} \Leftrightarrow A$ is $G_{2}$-instanton on $M^{7}$.

Rmk: $d_{A}^{*} F_{A}=0$ i.e. $A$ is a Yang-Mills connection! Topological
information e.g.
$Y M(A)=\int F_{A} \wedge * F_{A}=-\int F_{A} \wedge F_{A} \wedge \varphi=-\left[p_{1}(M) \wedge \varphi\right]$

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On $\left(M^{7}, \varphi\right)$ we have $\Omega^{2} \cong \Omega_{7}^{2} \oplus \Omega_{14}^{2}$ since $\Omega^{2} \cong \mathfrak{s o}(7) \cong \mathbb{R}^{7} \oplus \mathfrak{g}_{2}$. Call $F_{A}$ a $G_{2}$-instanton if $F_{A} \in \Omega_{14}^{2}$ i.e. $*\left(F_{A} \wedge \varphi\right)=-F_{A} \Leftrightarrow F_{A} \wedge * \varphi=0$.


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## Origin of deformed connections: Physics, SYZ conjecture,...

Mirror symmetry: $(P, \omega, \Omega) \longleftrightarrow(\check{P}, \check{\omega}, \check{\Omega})$ : Differential geometric version: SYZ conjecture: In certain limits (large J/vol), get dual special Lagrangian torus fibration st flat $U(1)$ connection $\mathbb{T}^{3} \longleftrightarrow$ Point on $\overleftarrow{T}^{3}$

MMMS, LYZ: Graph of $s: B^{3} \rightarrow P$ ' is special Lagrangian iff Connection on $P$ is dHYM i.e. $F_{A} \wedge \operatorname{Im}(\Omega)=0$ and $F_{A} \wedge \frac{1}{2} \omega \wedge \omega=\frac{1}{6} F_{A}^{3}$ (phase 1). Generally: $\operatorname{Im}\left(\left(\omega+i F_{A}\right)^{3}\right)=\tan (\theta) \operatorname{Re}\left(\left(\omega+i F_{A}\right)^{3}\right)$ since $\operatorname{Re}\left(e^{i \theta} \Omega\right)$ is a calibration

GYZ, LL: Replace "SLag $\mathbb{T}^{3}$ " by "co-associative $\mathbb{T}^{4}$ ". Then graph of $s: B^{3} \rightarrow \check{M}$ is associative ( +flat connection on $B^{3}$ ) iff Connection on $M$ is deformed $G_{2}: F_{A} \wedge * \varphi=\frac{1}{6} F_{A}^{3}$.

Note: In the non-deformed case " $F_{A}^{3 "}$ term is zero!

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MMMS, LYZ: Graph of $s: B^{3} \rightarrow \check{P}$ is special Lagrangian iff Connection on $P$ is dHYM i.e. $F_{A} \wedge \operatorname{Im}(\Omega)=0$ and $F_{A} \wedge \frac{1}{2} \omega \wedge \omega=\frac{1}{6} F_{A}^{3}$ (phase 1). Generally: $\operatorname{Im}\left(\left(\omega+i F_{A}\right)^{3}\right)=\tan (\theta) \operatorname{Re}\left(\left(\omega+i F_{A}\right)^{3}\right)$ since $\operatorname{Re}\left(e^{i \theta} \Omega\right)$ is a calibration


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## Trivial examples of $\mathrm{d} G_{2}$ instantons

Trivial here means arising by pullback from lower dimensions!
Prop. On $M^{7}=S^{1} \times P^{6}$ as before, $A$ is dHYM (phase 1) on $P^{6} \Leftrightarrow A$ is $\mathrm{d} G_{2}$-instanton on $M^{7}$. Likewise: If $M^{4}$ is HK 4-manifold and take Riemannian product with $\mathbb{T}^{3}$. Then $A$ is ASD-instanton iff $d G_{2}$ (and also $G_{2}$ ) on $M^{4} \times \mathbb{T}^{3}$

By contrast to $\mathrm{d} G_{2}$, lots of dHYM examples are known, we have uniqueness and also existence results in many cases! Existence of solutions links to a notion of stability of $L$ [Chen 20] (this verifies a mirror symmetry conjecture of Thomas-Yau). Crucial difference between $\mathrm{d} G_{2}$ and dHYM is that there is $d d^{c}$-lemma on Kähler manifolds! dHYM equation can be expressed as a Monge-Ampère type equation $F_{A}=F_{A_{0}}+d^{c} f$ ! No such thing in $G_{2}$-geometry

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## What do we know about deformed $G_{2}$-instantons?

(1) The $\mathrm{d} G_{2}$-equation arises also the critical point of a Chern-Simons type functional: Let $\mathbf{A}:=A_{0}+t\left(A-A_{0}\right)$ on $M^{7} \times[0,1]_{t}$ and consider

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\mathcal{F}(\mathbf{A}):=\frac{1}{2} \int_{M \times[0,1]} \frac{1}{12} F_{\mathbf{A}}^{4}-F_{\mathbf{A}}^{2} \wedge * \varphi .
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$A$ is critical point of $\mathcal{F}$ iff it is $d G_{2}$. [Karigiannis-Leung 07]
(2) The moduli space of $d G_{2}$-instantons on $\left(M^{7}, \varphi\right)$ compact has expected dimension 0 i.e. they are discrete. [Kawai-Yamamoto 20]
(3) The only known non-trivial examples are on nearly parallel $G_{2}$-manifolds, e.g. $S U(3) / U(1)$ and $S^{7}$, i.e. $\nabla \varphi \neq 0$ but instead $d \varphi=4 * \varphi$ [Lotay-Oliveira 20]. They can be used to distinguish between isometric $G_{2}$-structures!
(9) $\exists$ volume $V(A)=\int \sqrt{\operatorname{det}\left(I d+F_{A}^{\sharp}\right) \text { vol } \Rightarrow \text { Gradient flow: Line bundle }}$ MCF [Jacob-Yau 14, Kawai-Yamamoto 21]. Critical points contain $\mathrm{d} G_{2}$ and $\mathrm{d} G_{2}$ are global minimisers. (only need $\mathrm{cpt}+d \varphi=0$ !) Despite all this, there are no non-trivial examples on a $G_{2}$ manifold!

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## Examples of $G_{2}$ manifolds

The possibility of $G_{2}$ metrics was first suggested by work of Berger 1950s
(1) First local (non-trivial) examples: cones over NK [Bryant 87]
(2) Complete examples: $\mathbb{R}^{4} \times S^{3}, \Lambda_{-}^{2}\left(S^{4}\right), \Lambda_{-}^{2}\left(\mathbb{C P}^{2}\right)$ [Bryant-Salamon 89]
(3) Compact examples [Joyce, Kovalev, CHNP, J-K,..]
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Too hard to find examples on compact manifold. Try non-compact with
maximal possible symmetry i.e. cohomogeneity one. Why? PDEs become ODEs!

Coho 1 implies that $M / G$ is 1 dimensional i.e. $[0,1], S^{1},[0, \infty),(0,1)$
Ricci flat manifold + coho $1 G$ action implies only possibility is $[0, \infty)$ Let's focus on Bryant-Salamon manifolds!

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## Invariant connections on coho 1 manifolds

$M \simeq G / H_{2} \cup G / H_{1} \times \mathbb{R}^{+}$, where $H_{1} \subset H_{2} \subset G$ and $H_{2} / H_{1}$ is a sphere.
(1) $\$\left(S^{3}\right) \cong \mathbb{R}^{4}$
$S^{3} \simeq \frac{S U(2)^{2}}{S U(2)}$$\Lambda_{-}^{2}\left(S^{4}\right) \simeq \frac{S p(2)}{S p(1) S p(1)} \cup \frac{S p(2)}{U(1) S p(1)} \times \mathbb{R}^{+}$
(3) $\Lambda_{-}^{2}\left(\mathbb{C P}^{2}\right) \simeq \frac{S U(3)}{U(2)} \cup \frac{S U(3)}{T^{2}} \times \mathbb{R}^{+}$

Find connections on principal orbit $G / H_{1}$. Given gp homomorphism $\lambda: H_{1} \rightarrow K$ we define $K \rightarrow G \times_{\lambda} K \rightarrow G / H_{1}$. Canonical connection given by $d \lambda: \mathfrak{h}_{1} \rightarrow \mathfrak{k}$ and all other connections are given by $\Lambda:(\mathfrak{m}, \operatorname{Ad}) \rightarrow(\mathfrak{k}, \operatorname{Ad} \circ \lambda)$ $H$-morphism and $\mathfrak{g}=\mathfrak{h}_{1} \oplus \mathfrak{m}$. When $K=U(1)$ i.e. $\mathfrak{k} \cong \mathbb{R}$, then
(1) $\Lambda: 6 \mathbb{R} \rightarrow \mathbb{R}$
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Conclusion: Only in the first case we can get families of $U(1)$ connections on principal orbit! Look for $\mathrm{d}_{2}$-instantons on $\mathbb{R}^{4} \times S^{3}$ with
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## Extending solutions to singular orbit

Extending solutions from $G / H_{1} \times \mathbb{R}^{+}$to $G / H_{2}$. Need bundle extension, in the case of $S^{3}$ all bundles are trivial! [Eschenburg-Wang 00] gives condition for smooth extension at singular orbit. Get singular initial value problem!

A $G$-invariant tensor $T \in C^{\infty}\left(\otimes^{p} T M \otimes^{q} T^{*} M\right)$ can be identified with a 1-parameter family of $\mathrm{H}_{2}$-equivariant map


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A $G$-invariant tensor $T \in C^{\infty}\left(\otimes^{p} T M \otimes^{q} T^{*} M\right)$ can be identified with a 1-parameter family of $\mathrm{H}_{2}$-equivariant map

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T_{t}: V \supset S^{k} \rightarrow \otimes^{p}(V \oplus \mathfrak{p}) \otimes^{q}(V \oplus \mathfrak{p})
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where $\mathfrak{g}=\mathfrak{h}_{2} \oplus \mathfrak{p}$ and $S^{k}=H_{2} / H_{1}$ (working in normal bundle $V$ of singular orbit!)
Each $T_{t}$ is a finite sum of $H_{1}$-invariant tensors in $\otimes^{p}(V \oplus \mathfrak{p}) \otimes^{q}(V \oplus \mathfrak{p})$. For smoothness the coefficient functions must be even or odd depending on degree of the invariant tensor and lowest order term must be at least equal to degree of tensor.

## Abelian $G_{2}$-instantons

Smoothness condition for Bryant-Salamon and BGGG was worked out by Lotay-Oliveira. Using this they show

## Lotay-Oliveira 16

For Bryant-Salamon $G_{2}$-instantons are $A=\frac{r^{3}-c}{r} \sum_{i=1}^{3} c_{i} \eta_{i}$ they are $A=\frac{r^{2}-81 / 16}{r^{2}-9 / 16} c_{1} \eta_{1}+\frac{(r-9 / 4) e^{r}}{\sqrt{r}(r+9 / 4)^{2}}\left(c_{2} \eta_{2}+c_{3} \eta_{3}\right)$, where $r$ is distance to zero section, $\eta_{i}$ are vertical 1 -forms and $c_{i} \in \mathbb{R}$.

Rmk: Abelian $G_{2}$-instanton is a linear PDE, so ODE problem is much easier. Lotay-Oliveira also construct gauge $S U(2) G_{2}$-instantons (which is non-linear)! Note that $d G_{2}$ is always non-linear although gauge is $U(1)$

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## F 22

For the Bryant-Salamon cone $\mathrm{d}_{2}$-instantons are given by $A=f(r) \sum_{i=1}^{3} c_{i} \eta_{i}$, where $\log (c f(r)) f(r)^{2}=\frac{r^{4}}{\left(c_{1}^{2}+c_{2}^{2}+c_{3}^{2}\right)}: f \sim r^{2-\varepsilon}$. For BGGG, a smooth $\mathrm{d}_{2}$-instanton is given by $A=f(r) \eta_{1}$, where $24 \tan (f(r) / 3+c) f(r)=16 r^{2}-81: f \sim O(1)$. Here $c \in(0, \pi / 2)$

## Relation betweeen $G_{2}$ and $\mathrm{d} G_{2}$

Recall d $G_{2}$ means $F_{A} \wedge * \varphi=\frac{1}{6} F_{A}^{3}$. Note that to lowest order $F_{A}$ is a $G_{2}$-instanton: put differently in the "large volume limit" $d G_{2}$ is $G_{2}$ ! We can make this precise:

## Prop

Given $\left\{A_{k}\right\} d G_{2}$ such that $c_{k}:=\left\|F_{A_{k}}\right\|_{\infty} \rightarrow 0$ then define $B_{k}=A_{k} / c_{k}$ so that $\left\|F_{B_{k}}\right\|=1$, the limit $B_{k} \rightarrow B_{\infty}$ is a smooth $G_{2}$-instanton.

## Corollary <br> As $c \rightarrow \frac{\pi}{2}$, in BGGG case $A_{c}$ convergence to flat connection. Rescaling and

 taking limit yields the example of Lotay-Oliveira.One can construct local examples by evolving half-flat $S U(3)$-structures on nilmanifolds. [Chiossi-Salamon 02]
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## What about dSpin(7)?

Let $\left(M^{8}, \Phi\right)$ be a $\operatorname{Spin}(7)$ manifold, where $\Phi$ is the defining 4-form satisfying

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\left.\Phi\right|_{p}=d x_{0} \wedge \varphi+* \varphi
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in the analogous pointwise model as before. Similar properties as in $G_{2}$-case ( $\nabla \Phi=0$ iff $d \Phi=0$ : note that $\Phi=* \Phi$ ).
$\Phi$ is a calibration and calibrated submanifolds are called Cayley 4-folds. We have the irreducible splitting $\Omega^{2}(M)=\Omega_{7}^{2} \oplus \Omega_{21}^{2} \cong \mathbb{R}^{7} \oplus \operatorname{spin}(7)$. A connection $A$ is called $\operatorname{Spin}(7)$-instanton if $*\left(F_{A} \wedge \Phi\right)=-F_{A}$ i.e. $\pi_{7}^{2}\left(F_{A}\right)=0$ and a dSpin(7)-instanton if $\pi_{7}^{2}\left(F_{A}-\frac{1}{6} * F_{A}^{3}\right)=0$ and Note: If $F_{A}$ is $\operatorname{Spin}(7)$-instanton then $\pi_{7}^{4}\left(F_{A} \wedge F_{A}\right)=0$ since $S^{2}\left(\Omega_{21}^{2}\right)$ has no 7-dimensional component!
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## Example on a cone over $S^{7}$

As before, let's try to look for cohomogeneity one examples. Consider the spinor bundle of $S^{4}: \$\left(S^{4}\right)=\frac{S p(2)}{S p(1) S p(1)} \cup \frac{S p(2)}{S p(1)} \times \mathbb{R}^{+}$with Bryant-Salamon Spin(7) metric.


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The connection $A=f(r) \sum_{i=1}^{3} c_{i} \eta_{i}$

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f(r)=\frac{3 r^{2}}{10} W\left(c r^{128 / 27}\right)^{-1 / 2}: f \sim r^{2-\varepsilon}
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where $\eta_{i}$ are vertical 1-forms on $S^{3}$, gives an explicit solution on the Bryant-Salamon cone. $W$ is the Lambert $W$ function

Note: this is the cone over squashed $S^{7}$ not round one $(\mathrm{d} \operatorname{Spin}(7)$ is $\operatorname{Spin}(7)$ in the latter case with $S p(2) S p(1)$-symmetry and solution is $\left.f=c r^{2}\right)$ ! This applies to any squashed 3-Sasakian e.g. $S U(3) / U(1)$ !

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## dSpin(7) vs dHYM

As before we can construct trivial examples by taking products: $S^{1} \times G_{2}$, $\mathbb{T}^{2} \times C Y^{3}$ ! But there are also other possibilities!
In dimension 8, we have $S p(2) \subset S U(4) \subset \operatorname{Spin}(7)$ i.e. hyperKähler $\subset$ Calabi-Yau c $\operatorname{SPin}(7)$.


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Suppose $A$ is holomorphic wrt $(\omega, \Omega)$. Then dHYM wrt $\omega$ (phase 1) iff dSpin(7) wrt $\Phi$.


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## Work in progress...

Since $\mathfrak{s u}(3)=\mathfrak{u}(1) \oplus 3 \mathbb{R} \oplus \mathbb{R}^{2} \oplus \mathbb{R}^{2}$, on $\left(T^{*} \mathbb{C P}^{2}, \omega_{1}, \omega_{2}, \omega_{3}\right)$ we can express the metric by

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where $r \in[\sqrt{2 c}, \infty)$.
We have a 3-parameter family of connections

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and extension to $\mathbb{C P}^{2}$ depends on $H^{2}\left(\mathbb{C P}^{2}, \mathbb{Z}\right)=\mathbb{Z}$ and extension of line bundle: $a_{2}(\sqrt{2 c})=k, a_{3}(\sqrt{2 c})=a_{4}(\sqrt{2 c})=0 \Rightarrow F_{A(\sqrt{2 c})}=k \omega_{\mathbb{C P}^{2}}$.

There are 3 independent Kähler forms $\omega_{1}, \omega_{2}, \omega_{3}$ but the vector field $X_{1}$ (dual to $\theta_{1}$ ) permutes $\omega_{2}$ and $\omega_{3}$ ! Suffices to restrict to $\omega_{1}, \omega_{2}, \Phi_{1}, \Phi_{2}$

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\tan (\theta)=\frac{2\left(a_{2} r^{2}-2 c k\right)\left(a_{2}^{2}-r^{4} / 4+c^{2}-k^{2}\right)}{\left(a_{2}^{2}-r^{4} / 4+c^{2}-k^{2}\right)^{2}-\left(a_{2} r^{2}-2 c k\right)^{2}} .
$$

Rmk: $a_{2}=2 c k / r^{2}$ is $\operatorname{Sp}(2)$ connection [Hit] hence $\operatorname{Spin}(7)$ instanton wrt $\Phi_{2}$.
As comparison: You can play similar game on $F_{1,2}=S U(3) / \mathbb{T}^{2}$ with Kähler Einstein structure to get

$$
\tan (\theta)=\frac{a_{2}\left(a_{2}^{2}-k^{2}-3\right)}{3 a_{2}^{2}-k^{2}-1}
$$

note that here $a_{2}$ is a fixed constant! Rmk: $\theta$ is fixed by $a_{2}$ and $k$ because of compact setting!

## Results in progress...

## Thm...

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## Thm...

$A$ is dHYM wrt $\omega_{2}$ iff $a_{4}=0$ and $a_{2}(r)=2 c k / r^{2}$ is defined by

$$
\tan (\theta)=\frac{2\left(4 a_{3} r^{4} \sqrt{r^{4}-4 c^{2}}\right)\left(r^{8}-4 r^{4}\left(a_{3}^{2}+c^{2}-k^{2}\right)-16 c^{2} k^{2}\right)}{\left(4 a_{3} r^{4} \sqrt{r^{4}-4 c^{2}}\right)^{2}-\left(r^{8}-4 r^{4}\left(a_{3}^{2}+c^{2}-k^{2}\right)-16 c^{2} k^{2}\right)^{2}} .
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## Results in progress...

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$A$ is dHYM wrt $\Phi_{1}=\frac{1}{2}\left(-\omega_{1}^{2}+\omega_{2}^{2}+\omega_{3}^{2}\right)$ iff either of the following holds:

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A=3 k \theta_{1}+\frac{6 c k}{r^{2}} \theta_{2},
$$

or,

$$
A=3 k \theta_{1}+\frac{6 c k}{r^{2}} \theta_{2}+\frac{\sqrt{\left(r^{4}+36 k^{2}\right)\left(r^{4}-4 c^{2}\right)}}{2 r^{2} \sqrt{C_{3}^{2}+C_{4}^{2}}}\left(C_{3} \theta_{3}+C_{4} \theta_{4}\right)
$$

or,

$$
A=3 k \theta_{1}+\left(\frac{1}{2} C_{0} r^{2} \pm \frac{1}{2} \sqrt{C_{0}^{2} r^{4}-24 C_{0} c k+r^{4}-4 c^{2}+36 k^{2}}\right) \theta_{2} .
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## Rmk: Some of these examples are in fact dHYM for instance wrt $\omega_{2}$ (with

 phase 1: $C_{3}=0$ ). But more surprisingly this includes dHYM for all $\theta$ wrt $\omega_{1}$. I have not yet simplified the solutions wrt $\Phi_{2}$ :/
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## Merci beaucoup pour votre attention!


[^0]:    There are 3 independent Kähler forms $\omega_{1}, \omega_{2}, \omega_{3}$ but the vector field $X_{1}$ (dual to $\theta_{1}$ ) permutes $\omega_{2}$ and $\omega_{3}$ ! Suffices to restrict to $\omega_{1}, \omega_{2}, \Phi_{1}, \Phi_{2}$

